

The Twistor Transform of a Verlinde formula

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Introduction

Let Σ be a compact Riemann surface of genus g . The moduli space $\mathcal{M}_g = \mathcal{M}_g(2, 1)$ of stable rank 2 holomorphic bundles over Σ with fixed determinant bundle of degree 1 is a smooth complex $(3g-3)$ -dimensional manifold [25]. The anticanonical bundle of M is the square of a holomorphic line bundle L , some power of which embeds \mathcal{M}_g into a projective space. The dimensions of the vector spaces $H^0(\mathcal{M}_g, \mathcal{O}(L^{m-1}))$ of holomorphic sections of powers of L are known to be independent of the choice of complex structure on Σ , and are given by the formula

$$h^0(\mathcal{M}_g, \mathcal{O}(L^{m-1})) = -m^{g-1} \sum_{i=1}^{2m-1} (-1)^i \operatorname{cosec}^{2g-2}\left(\frac{i\pi}{2m}\right) \quad (0.1)$$

predicted by Verlinde [29]. This is closely related to the structure of the cohomology ring of \mathcal{M}_g , and a number of independent proofs and generalizations of (0.1) are now known. Below we shall follow closely the approach of Szenes [27].

In the case in which Σ is a hyperelliptic surface, and is therefore a 2-fold branched covering of \mathbb{CP}^k , Desale and Ramanan [9] exhibit \mathcal{M}_g as a complex submanifold of the flag manifold $\mathcal{F}_g = SO(2g+2)/(U(g-1) \times SO(4))$. As explained in [27] this reduces verification of (0.1) to certain $SO(2g+2)$ -equivariant calculations. Our contribution is to observe that \mathcal{F}_g is the twistor space of the real oriented Grassmannian $\mathcal{G}_g = SO(2g+2)/(SO(2g-2) \times SO(4))$ in the sense of [7, 8] for all $g \geq 3$. This enables us to relate the cohomology of the symmetric space \mathcal{G}_g directly to the cohomology of \mathcal{M}_g , and we obtain a set of generators for the latter which may be compared to the universal ones described in [21, 28, 10]. As a feasibility study, we illustrate the theory in the present paper for the case $g = 3$ which is worthy of special attention since the fibration $\mathcal{F}_3 \rightarrow \mathcal{G}_3$ encapsulates the quaternionic structure of the base space in a manner first identified by Wolf [31].

In the first section we investigate the cohomology of $\mathcal{G}_3 = SO(8)/(SO(4) \times SO(4))$. Using its quaternionic spin structure, we prove that the odd Pontrjagin classes of \mathcal{G}_3 vanish, and that its \hat{A} class simplifies remarkably. In the second section we recover Ramanan's description [23] of the Chern ring of \mathcal{M}_3 in the context of the natural mapping $\mathcal{M}_3 \rightarrow \mathcal{G}_3$, enabling $h^0(\mathcal{M}_3, L^{m-1})$ to be computed rapidly. Whilst this provides only a particularly simple instance of (0.1), results of the third section identify $H^0(\mathcal{M}_3, L^k)$ with a virtual representation of $SO(8)$ that also arises from the kernels of coupled Dirac operators on \mathcal{G}_3 . Similar techniques can in theory be applied to higher genus cases, and formulae such as $p_1^g = 0$ on \mathcal{M}_g [28, 16] may be expected to interact with properties of \mathcal{G}_g such as the constancy of the elliptic genera considered in [30, 15].

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1. Grassmannian cohomology

From now on we denote by \mathcal{G} the Grassmannian

$$\mathcal{G}_3 = \frac{SO(8)}{SO(4) \times SO(4)} \quad (0.1)$$

that parametrizes real oriented 4-dimensional subspaces of \mathbb{R}^8 . Let W denote the tautological real rank 4 vector bundle over \mathcal{G} , and W^\perp its orthogonal complement in the trivial bundle over \mathcal{G} with fibre \mathbb{R}^8 . The bundles W and W^\perp arise from the standard representations of the two $SO(4)$ factors constituting the isotropy subgroup in (0.1), and it follows that

$$T\mathcal{G} \cong W \otimes W^\perp. \quad (0.2)$$

The $SO(8)$ -invariant Riemannian metric on \mathcal{G} determines an isomorphism $W \cong W^*$ of vector bundles.

The decomposition (0.2) may be refined by lifting the $SO(4)$ structure of W to $Spin(4) \cong SU(2) \times SU(2)$ on a suitable open dense subset \mathcal{G}' of \mathcal{G} . This procedure is one that is familiar from the study of Riemannian 4-manifolds, and

$$W_{\mathbb{C}} \cong U \otimes_{\mathbb{C}} V,$$

where U and V are each complex rank 2 vector bundles over \mathcal{G}' . The resulting isomorphism

$$(T\mathcal{G})_{\mathbb{C}} \cong U \otimes (V \otimes W_{\mathbb{C}}^\perp) \quad (0.3)$$

reflects the fact that \mathcal{G} is a quaternion-Kähler manifold [31, 24]. In (0.3), U may be thought of as a quaternionic line bundle (usually called H), and its cofactor $V \otimes W_{\mathbb{C}}^\perp$ (usually called E) has structure group $SU(2) \times SO(4)$ extending to $Sp(4)$.

The Betti numbers of a quaternion-Kähler $4n$ -manifold of positive scalar curvature satisfy $b_{2k+1} = 0$ for all k and $b_{2k-4} \leq b_{2k}$ for $k \leq n+1$. They are also subject to the linear constraint of [17] which for $n = 4$ takes the form

$$3(b_2 + b_4) = 1 + b_6 + 2b_8.$$

This is well illustrated by \mathcal{G} , which has Poincaré polynomial

$$P_t(\mathcal{G}) = 1 + 3t^4 + 4t^8 + 3t^{12} + t^{16},$$

and is the only real Grassmannian to have $b_4 > 2$. (These facts may be deduced from [12, chapter XI].) We shall in fact only be concerned with the subring generated by the Euler class $e = e(W)$ and the first Pontrjagin class $f = p_1(W)$.

Although the classes e and f are very natural, it will ultimately be more convenient to consider

$$u = -c_2(U), \quad v = -c_2(V).$$

Because of the \mathbb{Z}_2 -ambiguity in the definition of U, V , the classes u, v are not integral, but the symmetric products $\odot^2 U, \odot^2 V$ are globally defined so $4u, 4v$ belong to $H^4(\mathcal{G}, \mathbb{Z})$. If we write formally $4u = \ell^2$ then

$$\text{ch}(U) = e^{\ell/2} + e^{-\ell/2} = 2 + u + \frac{1}{12}u^2 + \frac{1}{360}u^3 + \frac{1}{20160}u^4. \quad (0.4)$$

The class ℓ is given geometrical significance by the splitting (0.2). An analogous expression to (0.4) holds for $\text{ch}(V)$, and from $\text{ch}(W_{\mathbb{C}}) = \text{ch}(U)\text{ch}(V)$, we obtain

$$\begin{aligned} e &= u - v, \\ f &= 2(u + v). \end{aligned} \tag{0.5}$$

We may add that $p_2(W) = c_4(W_{\mathbb{C}}) = (u - v)^2$ confirming the well-known relation

$$p_2(W) = e^2. \tag{0.6}$$

Moreover, the space $H^4(\mathcal{G}, \mathbb{Z})$ is generated by e, f together with $e(W^{\perp})$ [19].

1.1 Proposition *Evaluation on the fundamental cycle $[\mathcal{G}]$ yields*

$$\begin{aligned} e^4 &= 2 = e^2 f^2, \quad e^3 f = 0 = e f^3, \quad f^4 = 4; \\ u^4 &= \frac{21}{64} = v^4, \quad u^3 v = -\frac{7}{64} = uv^3, \quad u^2 v^2 = \frac{5}{64}. \end{aligned}$$

We shall deduce these Schubert-type relations from a description of the total Pontrjagin class and the \hat{A} class

$$\begin{aligned} P(T\mathcal{G}) &= 1 + P_1 + P_2 + P_3 + P_4, \\ \hat{A}(T\mathcal{G}) &= 1 + \hat{A}_1 + \hat{A}_2 + \hat{A}_3 + \hat{A}_4 \end{aligned}$$

of the tangent bundle (0.2) of \mathcal{G} . (Upper case P_i 's are used to prevent a future clash of notation.) The classes \hat{A}_i , $1 \leq i \leq 4$ are determined in terms of the P_i in the usual way [14], and

1.2 Proposition $P_1 = 0 = P_3$ and $\hat{A}(\mathcal{G}) = 1 - \frac{1}{240}f^2$.

Proof of both propositions. It is easy to check that, in the presence of (0.5), the two sets of equations of Proposition 1.1 are equivalent. The equalities $u^4 = v^4$ and $u^3 v = uv^3$ are immediate from the symmetry between U and V , and these are equivalent to $ef^3 = 0 = ef^3$. Using (0.6), we have

$$\text{ch}(W_{\mathbb{C}}) = 4 + f + \frac{1}{12}(-2e^2 + f^2) + \frac{1}{360}(-3e^2 f + f^3) + \frac{1}{20160}(2e^4 - 4e^2 f^2 + f^4). \tag{0.7}$$

From (0.2) and (0.7),

$$\begin{aligned} \text{ch}(T\mathcal{G})_{\mathbb{C}} &= (\text{ch } W_{\mathbb{C}})(8 - \text{ch } W_{\mathbb{C}}) \\ &= 16 - f^2 + \frac{1}{6}(2e^2 f - f^3) + \frac{1}{720}(-20e^4 + 32e^2 f^2 - 9f^4). \end{aligned} \tag{0.8}$$

In particular $P_1 = 0$, and so we also have

$$\text{ch}(T\mathcal{G})_{\mathbb{C}} = 16 - \frac{1}{6}P_2 + \frac{1}{120}P_3 + \frac{1}{10080}(P_2^2 - 2P_4). \tag{0.9}$$

Comparing (0.8) and (0.9) gives

$$P_2 = 6f^2, \quad P_3 = 20(2e^2 f - f^3), \quad P_4 = 140e^4 - 224e^2 f^2 + 81f^4. \tag{0.10}$$

The remainder of the proof is based on the following less obvious facts.

(i) \mathcal{G} is a spin manifold (see forward to (0.1)) carrying a metric of positive scalar curvature. Therefore its \hat{A} genus

$$\hat{A}_4 = \frac{1}{2^{16}3^45^27} (762P_1^4 - 1808P_1^2P_2 + 416P_2^2 + 1024P_1P_3 - 384P_4) \quad (0.11)$$

vanishes. Thus

$$0 = 416(6f)^2 - 384(140e^4 - 224e^2f^2 + 81f^4) = 5376(-10e^4 + 16e^2f^2 - 3f^4). \quad (0.12)$$

(ii) The dimension d of the isometry group of any quaternion-Kähler 16-manifold with positive scalar curvature is given by

$$d = 7 - \frac{8}{3}P_1u^3 + 64u^4$$

[24, page 170]. In the present case, $d = \dim SO(8) = 28$ and we obtain

$$21 = 64u^4 = \frac{1}{4}(16e^4 + 24e^2f^2 + f^4). \quad (0.13)$$

(iii) On any compact quaternion-Kähler $4n$ -manifold M with positive scalar curvature and $n > 2$, the index

$$\hat{A}(M, \odot^2 U) = \left\langle \text{ch}(\odot^2 U) \hat{A}, [M] \right\rangle,$$

vanishes; this is a consequence of [24, Corollary 6.7] which is explained in [17]. Given that

$$\text{ch}(\odot^2 U) = 3 + 4u + \frac{4}{3}u^2 + \frac{8}{45}u^3 + \frac{4}{315}u^4,$$

$$\hat{A} = 1 - \frac{1}{24}P_1 - \frac{1}{2^53^25}P_2 - \frac{1}{2^63^35^17}P_3 = 1 - \frac{1}{240}f^2 + \frac{1}{1008}(2e^2f - f^3),$$

and $u = (2e + f)/4$, it follows that

$$24e^4 - 26e^2f^2 + f^4 = 0. \quad (0.14)$$

Proposition 1.1 now follows from (0.12), (0.13), (0.14), and it only remains to prove that $P_3 = 0$. Because of the symmetry between W and W^\perp , it suffices to prove that $P_3e = 0 = P_3f$, but this follows from (0.10) and Proposition 1.1. QED

Remark. The vanishing of \hat{A}_4 and (0.12) above is in fact equivalent to the vanishing of the index $\hat{A}(M, T)$ of the Dirac operator coupled to the tangent bundle (see (0.2)), essentially the so-called Rarita-Schwinger operator. This index is known to be equivariantly constant on any spin manifold with S^1 action [30], and always vanishes in the homogeneous setting [15].

2. The flag manifold and moduli space

We denote by \mathcal{F} the complex 9-dimensional homogeneous space

$$\mathcal{F}_3 = \frac{SO(8)}{U(2) \times SO(4)} \quad (0.1)$$

that parametrizes complex 2-dimensional subspaces Π of \mathbb{C}^8 that are isotropic with respect to a standard $SO(8)$ -invariant bilinear form. It has a complex contact structure that was studied in [31] and exhibits it as the twistor space of \mathcal{G} in the sense of [24]. Projecting Π to a real 4-dimensional subspace of \mathbb{R}^8 determines an $SO(8)$ -equivariant mapping $\pi: \mathcal{F} \rightarrow \mathcal{G}$, and each fibre of π is isomorphic to $SO(4)/U(2)$ and defines a rational curve in the complex manifold \mathcal{F} .

From standard facts about twistor spaces [6, 24, 22], one knows that $\text{Pic}(\mathcal{F})$ is generated by a holomorphic line bundle L on \mathcal{F} such that

- (i) the restriction of L to each fibre $\pi^{-1}(x) \cong \mathbb{CP}^k$ equals $\mathcal{O}(2)$;
- (ii) L^5 is isomorphic to the anticanonical bundle κ^{-1} of \mathcal{F} .

The line bundle L admits a square root over an open set \mathcal{G}' of \mathcal{G} on which U and V are defined, there is a C^∞ isomorphism

$$\pi^*U \cong L^{1/2} \oplus L^{-1/2}. \quad (0.2)$$

Let ℓ denote the fundamental class $c_1(L)$ in $H^2(\mathcal{F}, \mathbb{Z})$. From the Leray-Hirsch theorem, there is an identity $(\ell/2)^2 + \pi^*c_2(U) = 0$ of real cohomology classes. In terms of integral classes, and omitting π^* ,

$$\ell^2 = 4u. \quad (0.3)$$

In the notation of the Introduction, let $\mathcal{M} = \mathcal{M}_3$. Szenes exhibits the latter as the zero set of a non-degenerate holomorphic section $s \in H^0(\mathcal{F}, \mathcal{O}(\sigma^*))$, where $\sigma = \odot^2\tau$ and τ denotes the tautological rank 2 complex vector bundle acquired from the embedding $\mathcal{F} \subset \text{Gr}_2(\mathbb{C}^8)$. (Such a section s corresponds to a quadratic form on \mathbb{C}^8 , but we shall not mention this again until the end of Section 3.) From the coset description (0.1), it follows that

$$\tau \cong L^{-1/2} \otimes \pi^*V; \quad (0.4)$$

the right-hand side is well defined on \mathcal{F} , even though the individual factors only make sense locally (for example on $\pi^{-1}(\mathcal{G}')$). Since $V \cong V^*$, we have $\sigma^* \cong L \otimes \pi^*\odot^2V$. The resulting holomorphic structure on $\pi^*\odot^2V$ coincides with that induced in a standard way from the fact that \odot^2V has a self-dual connection on the quaternion-Kähler manifold \mathcal{G} , in the sense of [18]. In particular, $\pi^*\odot^2V$ is trivial over each fibre $\pi^{-1} \cong \mathbb{CP}^k$. From now on we shall write \odot^2V in place of $\pi^*\odot^2V$, and often omit tensor product signs.

The cohomology classes ℓ, u, v may be pulled back from both \mathcal{G} and \mathcal{F} to \mathcal{M} , and we shall denote the resulting elements of $H^i(\mathcal{M}, \mathbb{R})$ by the same symbols.

2.1 Proposition *On \mathcal{M} , $3u^2 + 10uv + 3v^2 = 0$, and evaluation on $[\mathcal{M}]$ yields*

$$u^3 = \frac{7}{2} = -v^3, \quad uv^2 = \frac{3}{2} = -u^2v.$$

Proof. The submanifold \mathcal{M} of \mathcal{F} is Poincaré dual to the Euler class $c_3(\sigma^*)$, which is readily computed from the formula $\text{ch}(\sigma^*) = e^\ell \text{ch}(\odot^2V)$ (see (0.4)) and equals $4\ell(u - v)$. Then, for example,

$$\langle u^3, [\mathcal{M}] \rangle = \langle u^3 c_3(\sigma^*), [\mathcal{F}] \rangle = \langle 4\ell(u^4 - u^3v), [\mathcal{F}] \rangle = 8 \langle u^4 - u^3v, [\mathcal{G}] \rangle = \frac{7}{2},$$

the last equality from Proposition 1.1. The evaluation of u^2v , uv^2 and v^3 follows in exactly the same way.

Since $H^8(\mathcal{M}, \mathbb{R}) \cong H^4(\mathcal{M}, \mathbb{R})$ is 2-dimensional [20], there must be a non-trivial linear relation $au^2 + buv + cv^2 = 0$. The solution $(a = c)/b = 3/10$ can be found by multiplying the left-hand side by u and v in turn. QED

The next result gives an independent derivation of the characteristic ring in the context of the twistor fibration $\mathcal{F} \rightarrow \mathcal{G}$.

2.2 Proposition *The Chern and Pontrjagin classes of \mathcal{M} are given by*

$$\begin{aligned} c_1 = 2\ell, \quad c_2 = 4(3u + v), \quad c_3 = 8\ell u, \quad c_4 = -\frac{112}{3}uv, \quad c_5 = c_6 = 0; \\ p_1 = -8(u + v), \quad p_2 = \frac{3}{8}p_1^2, \quad p_3 = 0. \end{aligned}$$

Proof. It is known [24] that the fibration π gives a C^∞ splitting of the holomorphic tangent bundle of \mathcal{F} :

$$T^{1,0}\mathcal{F} \cong L \oplus L^{1/2}(V \otimes W_{\mathbb{C}}^\perp).$$

Combining this with the isomorphism

$$T^{1,0}\mathcal{F}|_{\mathcal{M}} \cong T^{1,0}\mathcal{M} \oplus (L \odot^2 V)|_{\mathcal{M}},$$

we obtain

$$\begin{aligned} \text{ch}(T^{1,0}\mathcal{M}) &= e^\ell + e^{\ell/2} \text{ch}(V W_{\mathbb{C}}^\perp) - e^\ell \text{ch}(\odot^2 V) \\ &= e^\ell (1 + e^{-\ell/2} \text{ch} V (8 - \text{ch} W_{\mathbb{C}}) - \text{ch}(\odot^2 V)). \end{aligned}$$

This yields the required expressions for c_1, c_2, c_3 . We also get $c_4 = 28(u + v)^2$ which reduces to $-112uv/3$ from Proposition 2.1. We next obtain $c_5 = -32\ell v(u + v)$, so that $c_5\ell = 0$ and the vanishing of c_5 follows from the fact that $H^2(\mathcal{M}, \mathbb{R})$ is 1-dimensional [20]. Finally, all these equalities combine to yield

$$c_6 = \frac{1}{3}(504u^3 + 2824u^2v + 1928uv^2 + 120v^3),$$

and Proposition 2.1 implies that $c_6 = 0$. The Pontrjagin classes p_i of \mathcal{M} are now determined from the Chern classes by the usual relations. QED

Remark. The cohomology ring and Chern classes of \mathcal{M} were computed in [23, Theorem 4], and comparison with that shows that

$$h = \ell, \quad \nu = \frac{1}{2}(3u + v).$$

In general, it is known that the total Pontrjagin class of \mathcal{M}_g equals $(1 + \frac{1}{2g-2}p_1)^{2g-2}$ [21]. Moreover, $p_1^g = 0$ [16, 28] and $c_i = 0$ if $i > 2g - 2$ [11].

The above enable the dimension d_k of $H^0(\mathcal{M}, \mathcal{O}(L^k))$ to be computed quickly. For this purpose it is convenient to set $k = m - 1$.

2.3 Theorem $d_{m-1} = \frac{1}{45}m^2(11 + 20m^2 + 14m^4).$

Proof. Given that $c_1(T^{1,0}\mathcal{F}) = 2\ell$, the Todd class $\text{td}(T^{1,0}\mathcal{M})$ of \mathcal{M} equals

$$e^\ell \hat{A}(T\mathcal{M}) = e^\ell \left[1 - \frac{1}{24}p_1 + \frac{1}{2^7 3^2 5}(7p_1^2 - 4p_2) \right].$$

Using Propositions 2.1, 2.2 and the Riemann-Roch theorem, we obtain

$$\begin{aligned} d_{m-1} &= \left\langle e^{m\ell} \left(1 + \frac{1}{3}(u+v) - \frac{11}{135}uv\right), [\mathcal{M}] \right\rangle \\ &= -\frac{22}{135}m^2u^2v + \frac{2}{9}m^4(u^3 + u^2v) + \frac{4}{45}m^6u^3, \end{aligned}$$

and the result follows. QED

3. Equivariant indexes

In this section, we begin by considering the Dirac operator over the Grassmannian \mathcal{G} . Recall from (0.3) that the quaternionic structure of \mathcal{G} is characterized by the vector bundles $H = U$ and $E \cong V W_{\mathbb{C}}$ (juxtaposition denotes tensor product). For $p \geq 4$, the exterior power $\bigwedge^p E$ contains a proper subbundle $\bigwedge_0^p E$ with the property that $\bigwedge^p E \cong \bigwedge_0^p E \oplus \bigwedge^{p-2} E$ and, as described in [4], the total spin bundle Δ of \mathcal{G} decomposes as $\Delta_+ \oplus \Delta_-$ where

$$\begin{aligned} \Delta_+ &\cong \odot^4 U \oplus \odot^2 U \bigwedge_0^2 E \oplus \bigwedge_0^4 E, \\ \Delta_- &\cong \odot^3 U E \oplus U \bigwedge_0^3 E. \end{aligned} \tag{0.1}$$

The fact that all the summands on the right-hand side are globally defined confirms that \mathcal{G} is spin, though we shall not in fact need the decompositions (0.1).

Now let X be any other complex vector bundle over \mathcal{G} . The choice of a connection on X allows one to extend the Dirac operator on \mathcal{G} to an elliptic operator

$$D_X: \Gamma(\Delta_+ X) \longrightarrow \Gamma(\Delta_- X).$$

The index of this coupled Dirac operator is by definition $\dim(\ker D_X) - \dim(\operatorname{coker} D_X)$. This extends to a homomorphism $K(\mathcal{G}) \rightarrow \mathbb{Z}$, so that the index of D_X is also defined when X is a virtual vector bundle. The Atiyah-Singer index theorem [3] asserts that the index of D_X equals

$$\hat{A}(\mathcal{G}, X) = \left\langle \operatorname{ch}(X) \hat{A}(T\mathcal{G}), [\mathcal{G}] \right\rangle \tag{0.2}$$

In our situation, this fact is closely related to the Riemann-Roch theorem on \mathcal{F} which provides the following interpretation of d_k .

Theorem 3.1 *Let $X_k = \odot^{2k+4} U - \odot^{2k+2} U \odot^2 V + \odot^{2k} U \odot^2 V - \odot^{2k-2} U$, $k \geq 1$. Then $d_k = \hat{A}(\mathcal{G}, X_k)$.*

Proof. Let σ denote the rank 3 vector bundle $\odot^2 \tau$ as above, and let (k) denote the operation of tensoring with L^k . The description of \mathcal{M} as the zero set of a section of $\sigma^* \cong \odot^2 V(1)$ provides a Koszul complex

$$0 \rightarrow \mathcal{O}_{\mathcal{F}}(\bigwedge^3 \sigma(k)) \rightarrow \mathcal{O}_{\mathcal{F}}(\bigwedge^2 \sigma(k)) \rightarrow \mathcal{O}_{\mathcal{F}}(\sigma(k)) \rightarrow \mathcal{O}_{\mathcal{F}}(k) \rightarrow \mathcal{O}_{\mathcal{M}}(k) \rightarrow 0,$$

or equivalently,

$$0 \rightarrow \mathcal{O}_{\mathcal{F}}(k-3) \rightarrow \mathcal{O}_{\mathcal{F}}(\odot^2 V(k-2)) \rightarrow \mathcal{O}_{\mathcal{F}}(\odot^2 V(k-1)) \rightarrow \mathcal{O}_{\mathcal{F}}(k) \rightarrow \mathcal{O}_{\mathcal{M}}(k) \rightarrow 0.$$

It follows that

$$\chi(\mathcal{M}, \mathcal{O}(k)) = a_k - b_{k-1} + b_{k-2} - a_{k-3}, \quad (0.3)$$

where

$$a_k = \chi(\mathcal{F}, \mathcal{O}(k)), \quad b_k = \chi(\mathcal{F}, \mathcal{O}(\odot^2 V(k))). \quad (0.4)$$

These holomorphic Euler characteristics may be computed using the Riemann-Roch theorem and the cohomological version [24, 7.2] of the twistor transform; the result is

$$a_k = \hat{A}(\mathcal{G}, \odot^{2k+4} U), \quad b_k = \hat{A}(\mathcal{G}, \odot^{2k+4} U \odot^2 V). \quad (0.5)$$

Finally, Proposition 2.2 implies that the canonical bundle $\mathcal{K}(\mathcal{M})$ is isomorphic to L^{-2} , so by Serre duality and Kodaira vanishing, $H^i(\mathcal{M}, \mathcal{O}(k)) = 0$ for all $i \geq 1$ and $k \geq -1$. In particular, $\chi(\mathcal{M}, \mathcal{O}(k)) = \dim H^0(\mathcal{M}, \mathcal{O}(k))$ for all $k \geq -1$, and the theorem now follows from (0.3). QED

The isometry group $SO(8)$ of \mathcal{G} acts naturally on the cohomology groups over \mathcal{F} of the sheaves $\mathcal{O}(k)$, $\mathcal{O}(\odot^2 V(k))$ considered above. The integers a_k , b_k and

$$d_k = a_k - b_{k-1} + b_{k-2} - a_{k-3}$$

are therefore the dimensions of certain virtual $SO(8)$ -modules, and we identify these shortly.

Let $V(\gamma)$ denote the complex irreducible representation of $SO(8)$ with dominant weight γ , where $\gamma = (\lambda_1, \lambda_2, \lambda_3, \lambda_4)$ with $\lambda_1 \geq \lambda_2 \geq \lambda_3 \geq \lambda_4 \geq 0$. We adopt standard coordinates so that $V(1, 0, 0, 0) = \mathbb{C}^8$ is the fundamental representation, and $V(1, 1, 0, 0) = \mathfrak{so}(\mathfrak{g}, \mathbb{C})$ is the complexified adjoint representation.

3.2 Proposition *Let $A_k = V(k, k, 0, 0)$ and $B_k = V(k+1, k-1, 0, 0)$. Then $a_k = \dim A_k$ and $b_k = \dim B_k$.*

Proof. The Weyl dimension formula states that

$$\dim(V(\gamma)) = \prod_{\alpha \in R_+} \frac{\langle \alpha, d + \gamma \rangle}{\langle \alpha, d \rangle},$$

where R_+ denotes the set of positive roots and d is half of their sum. With the above coordinates,

$$R_+ = \{(1, 1, 0, 0), (1, 0, 1, 0), (1, 0, 0, 1), (0, 1, 1, 0), (0, 1, 0, 1), (0, 0, 1, 1), \\ (1, -1, 0, 0), (1, 0, -1, 0), (1, 0, 0, -1), (0, 1, -1, 0), (0, 1, 0, -1), (0, 0, 1, -1)\},$$

$d = (3, 2, 1, 0)$ and we obtain

$$\begin{aligned} \dim A_k &= \frac{1}{4320} (k+1)(k+2)^3 (2k+5)(k+3)^3 (k+4), \\ \dim B_k &= \frac{1}{1440} k(k+1)^2 (k+2)(2k+5)(k+3)(k+4)^2 (k+5). \end{aligned}$$

We claim that the right-hand sides are equal to a_k and b_k respectively. It follows from (0.4) that a_k and b_k are polynomials in k of degree 9, and by Serre duality,

$$a_{-k} = -a_{k-5}, \quad b_{-k} = -b_{k-5}, \quad k \in \mathbb{Z}. \quad (0.6)$$

By (0.6) and suitable vanishing theorems [5], $a_k = 0 = b_k$ for $k = -4, -3, -\frac{5}{2}, -2, -1$. In addition, \mathcal{F} has Todd genus $a_0 = 1 = -a_{-5}$, and $b_0 = 0 = b_{-5}$. Accordingly,

$$\begin{aligned} a_k &= \frac{1}{4320}(k+1)(k+2)(2k+5)(k+3)(k+4)\tilde{a}_k, \\ b_k &= \frac{1}{1440}k(k+1)(k+2)(2k+5)(k+3)(k+4)(k+5)\tilde{b}_k. \end{aligned}$$

where \tilde{a}_k is a quartic polynomial in k with $\tilde{a}_0 = 36$ and \tilde{b}_k is quadratic in k .

Let $n = 2k + 4$. The formulae (0.5) involve $\text{ch}(\odot^n U) = f(n)$, where

$$f(x) = \frac{e^{(x+1)\ell/2} - e^{-(x+1)\ell/2}}{e^{\ell/2} - e^{-\ell/2}} \quad (0.7)$$

(see (0.2)). To evade an explicit calculation of $\text{ch}(\odot^n U)$, we exploit the following formulae which are easily deduced from (0.7).

3.3 Lemma $f'(0) = \frac{\ell/2}{\tanh(\ell/2)}, \quad f''(0) = u.$

The right-hand side of the first equation is the series used in the definition of Hirzebruch's L-genus, and using (0.3) can be rewritten as

$$\begin{aligned} \left. \frac{d}{dn} \right|_{n=0} \text{ch}(\odot^n U) &= 1 - \sum_{j \geq 1} (-1)^j \frac{2^{2j} B_j}{(2j)!} u^{2j} \\ &= \frac{1}{2} \left(1 - \frac{1}{3}u - \frac{1}{45}u^2 + \frac{2}{945}u^3 - \frac{1}{4725}u^4 \right), \end{aligned}$$

where B_j are the Bernoulli numbers [14]. From above, we obtain

$$\left. \frac{d}{dk} \right|_{k=-2} a_k = \frac{1}{270}(u^4 + 2u^3v + u^2v^2) - \frac{1}{4725}u^4 = 0 = \left. \frac{d^2}{dk^2} \right|_{k=-2} a_k.$$

It follows that \tilde{a}_k is divisible by $(k+2)^2$, and by Serre duality by $(k+3)^2$. We obtain $\tilde{a}_k = (k+2)^2(k+3)^2$. The identification $\tilde{b}_k = (k+1)(k+4)$ is similar, and proceeds using a less-enlightening version of Lemma 3.3; we omit the details. QED

The following table displays some of the above dimension functions in terms of k .

k	0	1	2	3	4	5	6	7	8
a_k	1	28	300	1925	8918	32928	102816	282150	698775
b_k	0	35	567	4312	21840	85050	274890	772464	1945944
d_k	1	28	265	1392	5145	15100	37681	83392	168273

Applying Serre duality and Kodaira vanishing over \mathcal{F} , recalling that $\mathcal{K}(\mathcal{F}) \cong L^{-5}$, shows that there is in fact an $SO(8)$ -equivariant isomorphism $A_k \cong H^0(\mathcal{F}, \mathcal{O}(k))$. In particular, A_1 may be identified with both the space of holomorphic sections of L and the Lie algebra $\mathfrak{so}(8, \mathbb{C})$ of infinitesimal automorphisms of the contact structure of \mathcal{F} . There is an associated moment mapping $\mathcal{F} \rightarrow \mathbb{P}(\mathfrak{so}(8, \mathbb{C})^*) \cong \mathbb{CP}^{27}$ that identifies \mathcal{F} with the projectivization of the nilpotent orbit of minimal dimension [26]. Accordingly, the $SO(8)$ -equivariant linear mapping

$$\phi_k: \odot^k(H^0(\mathcal{F}, \mathcal{O}(1))) \longrightarrow H^0(\mathcal{F}, \mathcal{O}(k)) \quad (0.8)$$

is onto for all $k \geq 1$. Indeed, A_k is the irreducible summand of $\odot^k A_1$ of highest weight, and it suffices to show that the restriction of ϕ_k to A_k is an isomorphism. Observe that A_k contains a decomposable tensor product $\xi^{\otimes k}$ for some non-zero $\xi \in A_1$ and $\phi_k(\xi^{\otimes k})$, being the k th power of ξ regarded as a section of L , is also non-zero. The irreducibility of A_k and Schur's lemma establishes the claim.

A similar argument can be given to establish an $SO(8)$ -equivariant isomorphism $B_k \cong H^0(\mathcal{F}, \mathcal{O}(\odot^2 V(k)))$, given that $H^i(\mathcal{F}, \mathcal{O}(\odot^2 V(k)))$ vanishes for all $i > 0$ and $k \geq 0$. One considers the mapping

$$\psi_k: H^0(\mathcal{F}, \mathcal{O}(\odot^2 V(1))) \otimes H^0(\mathcal{F}, \mathcal{O}(k-1)) \longrightarrow H^0(\mathcal{F}, \mathcal{O}(\odot^2 V(k))),$$

in which $H^0(\mathcal{F}, \mathcal{O}(\odot^2 V(1)))$ is isomorphic to the irreducible 35-dimensional $SO(8)$ -module $\odot_0^2 \mathbb{C}^8$ with highest weight $(2, 0, 0, 0)$. The irreducible summand of highest weight in the tensor product is isomorphic to B_k and the restriction of ψ_k to this is an isomorphism.

The above arguments can be streamlined by applying more sophisticated twistor transform machinery contained, for example, in [5]. In particular, A_k and B_k are known to be isomorphic to the respective kernels of natural twistor operators

$$\begin{aligned} \alpha_k: \odot^{2k} U &\longrightarrow E \odot^{2k+1} U, \\ \beta_k: \odot^{2k} U \odot^2 V &\longrightarrow E \odot^{2k+1} U \odot^2 V. \end{aligned}$$

Recall that \mathcal{M} is the zero set of an element s of the space $B_1 \cong \odot_0^2 \mathbb{C}^8$. For suitable hyperelliptic surfaces Σ , the section s will be a real element; at each point of \mathcal{G} it then defines a section of $W \oplus W^\perp$, which is a trivial bundle with fibre \mathbb{R}^8 (see (0.1)). In these terms the element $\tilde{s} \in \ker \beta_1$ determined by s is essentially the image of s by the homomorphism

$$\odot^2(W \oplus W^\perp)_{\mathbb{C}} \rightarrow \odot^2 W_{\mathbb{C}} \rightarrow \odot^2 U \odot^2 V \cong \text{Hom}(\odot^2 V, \odot^2 U).$$

This may be used to describe \mathcal{M} as a ‘branched cover’ of a real subvariety of \mathcal{G} .

The Horrocks instanton bundle over $\mathbb{CP}^{\mathbb{A}}$ discussed at the end of [18] provides an analogous situation in which a geometric object is defined by a non-degenerate solution of a twistor equation over a homogeneous space. Such situations are worthy of more systematic investigation.

References

1. M.F. Atiyah, R. Bott: Yang-Mills equations over Riemann surfaces, *Phil. Trans. R. Soc. London* **308** (1982), 523–615.
2. M.F. Atiyah, N.J. Hitchin, I.M. Singer: Self-duality in four-dimensional Riemannian geometry, *Proc. Roy. Soc. Lond.* **A362** (1978), 425–461.
3. M.F. Atiyah, I.M. Singer: The index theory of elliptic operators III, *Ann. Math.* **87** (1968), 546–604.
4. R. Barker, S.M. Salamon: Analysis on a generalized Heisenberg group, *J. London Math. Soc.* **28** (1983), 184–192.
5. R.J. Baston and M.G. Eastwood: *The Penrose transform: its interaction with representation theory*, Oxford University Press (1989).
6. A. Besse: *Einstein manifolds*, Springer (1987).
7. R.L. Bryant: Lie groups and twistor spaces, *Duke Math. J.* **52** (1985), 223–261.
8. F.E. Burstall, J.H. Rawnsley: *Twistor theory for Riemannian symmetric spaces*, *Lect. Notes Math.* **1424**, Springer (1990).
9. U.V. Desale, S. Ramanan: Classification of vector bundles of rank 2 over hyperelliptic curves, *Invent. Math.* **38** (1976), 161–185.
10. S.K. Donaldson: Gluing techniques in the cohomology of moduli spaces, in *Topological Methods in Modern Mathematics*, Publish or Perish (1993), 137–170.
11. D. Gieseker: A degeneration of the moduli space of stable bundles, *J. Diff. Geometry* **19** (1984), 173–206.
12. W. Greub, S. Halperin, R. Vanstone: *Curvature, connections and characteristic classes*, Volume 3, Academic Press (1976).
13. G. Harder, M.S. Narasimhan: On the cohomology groups of moduli spaces of vector bundles on curves, *Math. Ann.* **212** (1978), 215–248.
14. F. Hirzebruch: *Topological methods in algebraic geometry*, 3rd edition, Springer (1966).
15. F. Hirzebruch, P. Sladowy: Elliptic genera, involutions, and homogeneous spin manifolds, *Geom. Dedicata* **35** (1990) 309–343.
16. F.C. Kirwan: The cohomology rings of moduli spaces of bundles over Riemann surfaces, *J. Amer. Math. Soc.* **5** (1992), 853–906.
17. C.R. LeBrun, S.M. Salamon: Strong rigidity of positive quaternion-Kähler manifolds, *Invent. Math.* **118** (1994), 109–132.
18. M. Mamone Capria, S. Salamon: Yang-Mills fields on quaternionic spaces, *Nonlinearity* **1** (1988), 517–530.
19. J.W. Milnor, J.D. Stasheff: *Characteristic classes*, *Annals of Math. Studies* 76, Princeton University Press (1974).

- 20.** P.E. Newstead: Topological properties of some spaces of stable bundles, *Topology* **6** (1967), 241–262.
- 21.** P.E. Newstead: Characteristic classes of stable bundles over an algebraic curve, *Trans. Am. Math. Soc.* **169** (1972), 337–345.
- 22.** Y.S. Poon, S.M. Salamon: Eight-dimensional quaternion-Kähler manifolds with positive scalar curvature, *J. Diff. Geometry* **33** (1991), 363–378.
- 23.** S. Ramanan: The moduli space of vector bundles over an algebraic curve, *Math. Ann.* **200** (1973), 69–84.
- 24.** S. Salamon: Quaternionic Kähler manifolds, *Invent. Math.* **67** (1982), 143–171.
- 25.** C.S. Seshadri: Space of unitary vector bundles on a compact Riemann surface, *Ann. of Math.* **85** (1967), 303–336.
- 26.** A.F. Swann: Hyperkähler and quaternionic Kähler geometry, *Math. Ann.* **289** (1991), 421–450.
- 27.** A. Szenes: Hilbert polynomials of moduli spaces of rank 2 vector bundles I, *Topology* **32** (1993), 587–597.
- 28.** M. Thaddeus: Conformal field theory and the moduli space of stable bundles, *J. Diff. Geometry* **35** (1992), 131–149.
- 29.** E. Verlinde: Fusion rules and modular transformations in 2d conformal field theory, *Nucl. Phys. B* **300** (1988), 360–376.
- 30.** E. Witten: The index of the Dirac operator in loop space, in *Elliptic curves and modular forms in algebraic topology*, *Lect. Notes Math.* **1326**, Springer (1988), 161–181.
- 31.** J.A. Wolf: Complex homogeneous contact structures and quaternionic symmetric spaces, *J. Math. Mech.* **14** (1965) 1033–1047.

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